

with periodic boundary conditions. This extends an earlier theorem of Dávila in [6] for $s \leq 2$.

Before stating the stabilization result mentioned above, we first need to ensure the well-posedness of the system. This was addressed by Dávila in [3] (see also [4]) under the following conditions on the coefficients:

$$(1.2) \quad \begin{aligned} a_3^2 b_2 &< 1 \text{ and } r = 0 \\ b_2 a_1 a_3 - b_1 a_3 + b_1 a_2 - a_2 &= 0 \\ b_1 a_1 - a_1 - b_1 a_2 a_3 + a_3 &= 0 \\ b_1 a_2^2 + b_2 a_1^2 - b_1 a_1 - a_2 &= 0. \end{aligned}$$

Indeed, under conditions (1.2), Dávila and Chaves [7] derived some conservation laws for the solutions of (1.1). Combined with an approach introduced in [2, 17], these conservation laws allow them to establish the global well-posedness in $H_p^s(0, 1)$, for any $s \geq 0$. Moreover, the authors also give a simpler derivation of the conservation laws discovered by Gear and Grimshaw, and Bona et al [1]. We also observe that these conservation properties were obtained employing the techniques developed in [13] for the single KdV equation; see also [12].

The well-posedness result reads as follows:

Theorem 1.1. *Assume that condition (1.2) holds. If $\phi, \psi \in H_p^s(0, 1)$ for some integer $s \geq 3$, then the system (1.1) has a unique solution satisfying*

$$u, v \in C([0, \infty); H_p^s(0, 1)) \cap C^1([0, \infty); H_p^{s-3}(0, 1)).$$

Moreover, the map $(\phi, \psi) \mapsto (u, v)$ is continuous from $(H_p^s(0, 1))^2$ into

$$(C([0, \infty); H_p^s(0, 1)) \cap C^1([0, \infty); H_p^{s-3}(0, 1)))^2.$$

For $k = 0$, the analogous theorem on the whole real line $-\infty < x < \infty$ was proved Bona et al. [1], for all $s \geq 1$.

With the global well-posedness result in hand, we can focus on the stabilization problem. For simplicity of notation we consider only the case

$$(1.3) \quad b_1 = b_2 = 1.$$

Then the conditions (1.2) take the simplified form

$$(1.4) \quad r = 0, \quad a_1^2 + a_2^2 = a_1 + a_2, \quad |a_3| < 1, \quad \text{and} \quad (a_1 - 1)a_3 = (a_2 - 1)a_3 = 0.$$

Hence either $a_3 = 0$ and $a_1^2 + a_2^2 = a_1 + a_2$, or $0 < |a_3| < 1$ and $a_1 = a_2 = 1$.

We prove the following theorem:

Theorem 1.2. *Assume (1.3) and (1.4). If $\phi, \psi \in H_p^s(0, 1)$ for some integer $s \geq 3$, then the solution of (1.1) satisfies the estimate*

$$\|u(t) - [u(t)]\|_{H_p^s(0,1)} + \|v(t) - [v(t)]\|_{H_p^s(0,1)} = o\left(e^{-k't}\right), \quad t \rightarrow \infty$$

for each $k' < k$.

An analogous theorem was proved in [10] for the usual KdV equation by using the infinite family of conservation laws for this equation. Such conservations lead to the construction of a suitable Lyapunov function that gives the exponential decay of the solutions. Here, we follow the same approach making use of the results established by Dávila and Chavez [7]. They proved that under the assumptions (1.2) system (1.1) also has an infinite family of conservation laws, and they conjectured the above theorem for this case. At this point we observe that some computations are simplified if we change u, v, ϕ and ψ to $u - [u], v - [v], \phi - [\phi]$ and $\psi - [\psi]$. Then, the new unknown functions u and v satisfy the same system (1.1) with ku and $instead$

of $k(u - [u])$ and $k(v - [v])$. Hence we consider the solutions of the simplified system

$$(1.5) \quad \begin{cases} u' + uu_x + u_{xxx} + a_3v_{xxx} + a_1vv_x + a_2(uv)_x + ku = 0, \\ v' + vv_x + v_{xxx} + a_3u_{xxx} + a_2uu_x + a_1(uv)_x + kv = 0, \\ u(0, x) = \phi(x), \\ v(0, x) = \psi(x) \end{cases}$$

with periodic boundary conditions, corresponding to initial data ϕ, ψ with zero mean values.

In order to obtain the result, we prove a number of identities and estimates for the solutions of (1.1). In view of Theorem 1.1 it suffices to establish these estimates for *smooth solutions*, i.e., to solutions corresponding to C^∞ initial data ϕ, ψ with periodic boundary conditions. For such solutions all formal manipulations in the sequel will be justified.

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INSTITUTO DE MATEMÁTICA, UNIVERSIDADE FEDERAL DO RIO DE JANEIRO, C.P. 68530 - CIDADE UNIVERSITÁRIA - ILHA DO FUNDÃO, 21945-970 RIO DE JANEIRO (RJ), BRAZIL
E-mail address: capistrano@im.ufrj.br

DÉPARTEMENT DE MATHÉMATIQUE, UNIVERSITÉ DE STRASBOURG, 7 RUE RENÉ DESCARTES, 67084 STRASBOURG CEDEX, FRANCE
E-mail address: vilmos.komornik@math.unistra.fr

INSTITUTO DE MATEMÁTICA, UNIVERSIDADE FEDERAL DO RIO DE JANEIRO, C.P. 68530 - CIDADE UNIVERSITÁRIA - ILHA DO FUNDÃO, 21945-970 RIO DE JANEIRO (RJ), BRAZIL
E-mail address: ademir@im.ufrj.br